

Calculus Rules

In standard, non-stochastic calculus, one computes a derivative or an integral using various rules. In the Itô stochastic calculus, one extends these rules to the stochastic terms.

Suppose that u is some function $u(x)$ of x . We want to express the differential du in terms of the differential dx .

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Taylor Series

Consider the Taylor series expansion of $u(x)$ about some value \bar{x} :

$$u(x) = u(\bar{x}) + u'(\bar{x})(x - \bar{x}) + \frac{1}{2}u''(\bar{x})(x - \bar{x})^2 + \frac{1}{3!}u'''(\bar{x})(x - \bar{x})^3 + \dots$$

Under certain general conditions, $u(x)$ equals this infinite sum exactly.

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Rewrite this expression in terms of the changes $\Delta x := x - \bar{x}$ and $\Delta u := u(x) - u(\bar{x})$:

$$\Delta u = u'(\bar{x})\Delta x + \frac{1}{2}u''(\bar{x})(\Delta x)^2 + \frac{1}{3!}u'''(\bar{x})(\Delta x)^3 + \dots$$

Replace the difference by the differential:

$$du = u'(\bar{x})dx + \frac{1}{2}u''(\bar{x})(dx)^2 + \frac{1}{3!}u'''(\bar{x})(dx)^3 + \dots \quad (1)$$

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Non-Stochastic Calculus

In standard, non-stochastic calculus, one computes a differential simply by keeping the first-order terms. For small changes in the variable, second-order and higher terms are negligible compared to the first-order terms. Equation (1) becomes

$$du = u' dx.$$

The change in u is proportional to the change in x .

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Stochastic Calculus—Itô's Formula

In stochastic calculus, one must also keep the second-order terms. Equation (1) becomes *Itô's formula*,

$$du = u' dx + \frac{1}{2}u''(dx)^2 \quad (2)$$

This equation is *exact*; the third-order and higher order terms are zero.

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Rules of Stochastic Calculus

One computes Itô's formula (2) using the rules (3). Let z denote Wiener-Brownian motion, and let t denote time. One computes using the rules

$$\begin{aligned} (dz)^2 &= dt, \\ dz dt &= 0, \\ (dt)^2 &= 0. \end{aligned} \quad (3)$$

The key rule is the first and is what sets stochastic calculus apart from non-stochastic calculus.

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Computation

Although we prove the rules (3) below, first let us consider the implication of the rules. One computes mechanically, as in ordinary algebra, but using the rules. The second-order terms cannot be dropped, since $(dz)^2 = dt$.

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Example

If $dx = m dt + s dz$, then

$$\begin{aligned}(dx)^2 &= (m dt + s dz)^2 \\ &= (m dt)^2 + (s dz)^2 + 2(m dt)(s dz) \\ &= 0 + s^2 dt + 0 \\ &= s^2 dt.\end{aligned}$$

The second-order term is non-zero, as long as the instantaneous stochastic part is non-zero ($s \neq 0$).

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Therefore Itô's formula (2) says

$$\begin{aligned}du &= u'(m dt + s dz) + \frac{1}{2}u''(m dt + s dz)^2 \\ &= u'(m dt + s dz) + \frac{1}{2}u''s^2 dt \\ &= \left(u'm + \frac{1}{2}u''s^2\right) dt + u's dz.\end{aligned}$$

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Third-Order and Higher-Order Terms

Like non-stochastic calculus, third-order and higher-order terms are zero. For example,

$$(dx)^3 = dx(dx)^2 = (m dt + s dz)s^2 dt = ms^2(dt)^2 + s dz dt = 0,$$

applying the rules.

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Square of Wiener-Brownian Motion

Consider $u = z^2$:

$$\begin{aligned}du &= u' dz + \frac{1}{2}u''(dz)^2 \\ &= 2z dz + \frac{1}{2}2(dz)^2 \\ &= 2z dz + dt.\end{aligned}$$

Relative to non-stochastic calculus, dt is an extra term.

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Confirmation of Previous Result

Essentially the same calculation confirms our earlier limiting result that $du = 2z dz$, with initial value $u(0) = 0$, has the solution $u = z^2 - t$:

$$\begin{aligned}du &= u_z dz + u_t dt + \frac{1}{2}u_{zz}(dz)^2 + u_{zt} dz dt + \frac{1}{2}u_{tt}(dt)^2 \\ &= 2z dz + (-1) dt + \frac{1}{2}2(dz)^2 + 0 dz dt + \frac{1}{2}0(dt)^2 \\ &= 2z dz - dt + dt \\ &= 2z dz.\end{aligned}$$

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Stochastic Exponential

If $u = e^{z-t/2}$, then

$$\begin{aligned} u_z &= u & u_{zz} &= u \\ u_t &= -\frac{1}{2}u & u_{zt} &= -\frac{1}{2}u & u_{tt} &= \frac{1}{4}u. \end{aligned}$$

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The Taylor series is

$$\begin{aligned} du &= u_z dz + u_t dt + \frac{1}{2}u_{zz} (dz)^2 + u_{zt} dz dt + \frac{1}{2}u_{tt} (dt)^2 \\ &= u dz - \frac{1}{2}u dt + \frac{1}{2}u (dz)^2 - \frac{1}{2}u dz dt + \frac{1}{2} \left(\frac{1}{4}u \right) (dt)^2 \\ &= u dz - \frac{1}{2}u dt + \frac{1}{2}u dt \\ &= u dz. \end{aligned}$$

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Logarithm

$$\begin{aligned} d \ln x &= \frac{d \ln x}{dx} dx + \frac{1}{2} \frac{d^2 \ln x}{dx^2} (dx)^2 \\ &= \left(\frac{1}{x} \right) dx + \frac{1}{2} \left(-\frac{1}{x^2} \right) (dx)^2 \\ &= \frac{dx}{x} - \frac{1}{2} \left(\frac{dx}{x} \right)^2. \end{aligned}$$

Hence the change $d \ln x$ in the logarithm is *not* the growth rate dx/x , unless the instantaneous stochastic part of dx is zero.

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Inverse

We have

$$(1 - dx)^{-1} = 1 + dx + (dx)^2,$$

as

$$(1 - dx) \left[1 + dx + (dx)^2 \right] = 1.$$

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Product Rule

$$\begin{aligned} d(xy) &= \frac{\partial(xy)}{\partial x} dx + \frac{\partial(xy)}{\partial y} dy \\ &\quad + \frac{1}{2} \frac{\partial^2(xy)}{\partial x^2} (dx)^2 + \frac{\partial^2(xy)}{\partial x \partial y} dx dy + \frac{1}{2} \frac{\partial^2(xy)}{\partial y^2} (dy)^2 \\ &= y dx + x dy + 0(dx)^2 + 1 dx dy + 0(dy)^2 \\ &= y dx + x dy + dx dy. \end{aligned}$$

Compared to non-stochastic calculus, $dx dy$ is an extra term.

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Error Rule

We prove the fundamental error rule $(dz)^2 = dt$, by taking the limit of the discrete-time analogue. Divide the time interval from zero to t into n periods of length Δt , so $t = n\Delta t$. Holding t fixed, define

$$\int_0^t (dz)^2 := \lim_{\Delta t \rightarrow 0} \sum_{i=1}^n [\Delta z_{(i-1)\Delta t}]^2.$$

Defining $e_i := \Delta z_{(i-1)\Delta t} = z_{i\Delta t} - z_{(i-1)\Delta t}$, we can restate this equation as

$$\int_0^t (dz)^2 := \lim_{n \rightarrow \infty} (e_1^2 + \dots + e_n^2).$$

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We rewrite the sum of the squared errors as

$$e_1^2 + e_2^2 + \dots + e_n^2 = t \left\{ \frac{1}{n} \left[\left(\frac{e_1^2}{\Delta t} \right) + \left(\frac{e_2^2}{\Delta t} \right) + \dots + \left(\frac{e_n^2}{\Delta t} \right) \right] \right\}.$$

Holding $t = n\Delta t$ fixed, take the limit as $\Delta t \rightarrow 0$, $n \rightarrow \infty$.

The expression in braces is the sample mean of n independent $\chi^2(1)$ variables. By the law of large numbers, the sample mean converges to the true mean 1 as the sample size increases.

Hence

$$\lim_{n \rightarrow \infty} (e_1^2 + e_2^2 + \dots + e_n^2) = t.$$

Therefore

$$\int_0^t (dz)^2 = t,$$

regardless of t . Of course

$$\int_0^t dt = t.$$

Comparing the two integrals proves

$$(dz)^2 = dt.$$

Time Rule

We next prove $(dt)^2 = 0$. Divide the time interval from zero to t into n periods of length Δt , so $t = n\Delta t$. By definition,

$$\begin{aligned} \int_0^t (dt)^2 &:= \lim_{\Delta t \rightarrow 0} \sum_{i=1}^n (\Delta t)^2 \\ &= \lim_{\Delta t \rightarrow 0} [n(\Delta t)^2] \\ &= \lim_{\Delta t \rightarrow 0} (n\Delta t) \lim_{\Delta t \rightarrow 0} \Delta t \\ &= t0 \\ &= 0, \end{aligned}$$

and the result follows.

Cross-Product Rule

The rule $dz dt = 0$ can be shown by a similar limiting argument.